

CP violation studies in $B \rightarrow D^{(*)}K^{(*)}$ in BaBar and Belle

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Abstract. BaBar and Belle results on $B \rightarrow D^{(*)}K^{(*)}$ decays are reviewed in the context of constraining the ϕ_3 angle in the Unitarity Triangle.

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1 Introduction

CP violation (CPV) studies in the B-sector are targeted at measurements of the three angles of the Unitarity Triangle. The angle $\phi_3(\gamma) \equiv \arg(V_{ud}V_{ub}^*/V_{cd}V_{cb}^*)$ can be observed through interference between $b \rightarrow c$ and $b \rightarrow u$ transitions. Good processes for such studies are strangeness-changing B decays to neutral charmed mesons:

$$B \rightarrow DX_s. \quad (1)$$

D stands for D^0 or \bar{D}^0 as well as their excited states and X_s denotes a system of light mesons with $S = \pm 1$. Example diagrams for such decays are shown in Fig. 1. The decays $B^- \rightarrow D^0K^-$ are driven by $b \rightarrow c$ transitions and $B^- \rightarrow \bar{D}^0K^-$ by $b \rightarrow u$ transitions¹. When D^0 and \bar{D}^0 decay to a common final state, the two amplitudes (denoted $A_{b \rightarrow c}$ and $A_{b \rightarrow u}$) interfere, leading to direct CPV. Processes of this type, being free from penguin contributions, provide a theoretically clean environment to extract ϕ_3 .

The first idea of measuring ϕ_3 in decays of the type (1) came from Gronau, London and Wyler (GLW) [1]. In this method interference between the two amplitudes should be observed in $B \rightarrow D_{\pm}K$ decay modes, where $D_{\pm} \equiv (D^0 \pm \bar{D}^0)/\sqrt{2}$ denotes D meson CP eigenstates. In such case ϕ_3 can be extracted from CP asymmetries and ratios of decay rates using the following relations [2]:

$$\begin{aligned} \mathcal{A}_{\pm} &\equiv \frac{\mathcal{B}(B^- \rightarrow D_{\pm}K^-) - \mathcal{B}(B^+ \rightarrow D_{\pm}K^+)}{\mathcal{B}(B^- \rightarrow D_{\pm}K^-) + \mathcal{B}(B^+ \rightarrow D_{\pm}K^+)} \\ &= \frac{\pm 2r \sin \delta \sin \phi_3}{1 + r^2 \pm 2r \cos \delta \cos \phi_3} \end{aligned} \quad (2)$$

$$\begin{aligned} \mathcal{R}_{\pm} &\equiv \frac{2(\mathcal{B}(B^- \rightarrow D_{\pm}K^-) + \mathcal{B}(B^+ \rightarrow D_{\pm}K^+))}{\mathcal{B}(B^- \rightarrow D^0K^-) + \mathcal{B}(B^+ \rightarrow \bar{D}^0K^+)} \\ &= 1 + r^2 \pm 2r \cos \delta \cos \phi_3. \end{aligned} \quad (3)$$

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¹ Charge conjugation is implied throughout the paper unless explicitly stated otherwise.

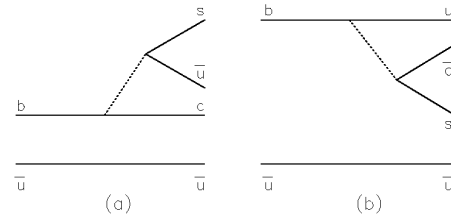


Fig. 1. Diagrams corresponding to $A_{b \rightarrow c}$ (a) and $A_{b \rightarrow u}$ (b) amplitudes

In the above equations δ is the relative strong phase of the interfering amplitudes and $r \equiv |A_{b \rightarrow u}/A_{b \rightarrow c}|$. The observables \mathcal{A}_{\pm} and \mathcal{R}_{\pm} contain information sufficient to extract δ and r , and thus ϕ_3 can be measured without hadronic uncertainties. Sensitivity of such measurements, however, depends on the ratio r which for decays with the biggest branching fractions tends to be small. (for $B^- \rightarrow DK^-$ the expected value of r is $\sim 0.1 \div 0.2$). This is mainly due to the color suppression of the $A_{b \rightarrow u}$ amplitude. Many subsequent approaches, which to a large extent can be considered as variants and extensions of the GLW method, try to overcome this difficulty. Altogether they offer a rich experimental program exploiting a variety of final states and observables in decays of the type (1).

2 Experimental results

While a large number of methods have been proposed, all of them present a challenge for experiments. Among the main problems one should mention are: small product branching fractions (e.g. D^0 decays to CP-eigenstates, f_{CP} , constitute $\sim 1\%$ of the total width leading to $\mathcal{B}(B^- \rightarrow D^0K^-) \times \mathcal{B}(D^0 \rightarrow f_{CP}) \sim 10^{-6}$); observables which weakly depend on interference effects or are difficult to measure and background from much more abundant non-strange modes like $B \rightarrow D\pi$. (e.g. one estimates $\frac{\mathcal{B}(B^- \rightarrow D^0K^-)}{\mathcal{B}(B^- \rightarrow D^0\pi^-)} \approx 0.075$ from the value of the Cabibbo angle).

Consequently, all the methods require a huge data sample (typically of the order of 500 M $B\bar{B}$'s or more) and high performance of a multipurpose detector. Experiments at B-factories are potentially a good place to perform such measurements. In particular they provide good kinematical resolution and K/π separation, which are essential for these studies. With the present luminosities the data sample of 500 fb^{-1} can be collected after a few years of running. The results reported here are based on data samples typically of $\approx 80 \text{ fb}^{-1}$ and represent preliminary studies towards the ϕ_3 measurement.

2.1 $B^- \rightarrow DK^{*-}$

$B^- \rightarrow D^0 K^-$ decays have been observed by experiments running at $\Upsilon(4S)$. Recent measurements of the ratios of branching fractions of Cabibbo-suppressed to Cabibbo-favored modes $R = \frac{\mathcal{B}(B^- \rightarrow D^0 K^-)}{\mathcal{B}(B^- \rightarrow D^0 \pi^-)}$ are summarized in Table 1. The results are consistent with the Cabibbo suppression mentioned above. Belle and BaBar have also studied $B^- \rightarrow D_{\pm} K^-$ modes with data samples of 78 fb^{-1} and 82 fb^{-1} respectively [4], [5]. Both experiments found signals in the channels with D decaying to CP-even eigenstates: $D_+ \rightarrow K^+ K^-, \pi^+ \pi^-$. Belle also observed CP-odd modes: $D_- \rightarrow K_s^0 \pi^0, K_s^0 \phi, K_s^0 \omega, K_s^0 \eta, K_s^0 \eta'$. Figures 2 and 3 present ΔE distributions for these channels (ΔE is a difference between reconstructed B-meson energy and its nominal value in the $\Upsilon(4S)$ rest frame).

Results on \mathcal{A}_{\pm} and \mathcal{R}_{\pm} are summarized in Table 2. Measurement errors preclude direct determination of ϕ_3 from these data but meaningful constraints should be feasible in the near future by means of inequalities resulting from (2, 3): $\sin^2 \phi_3 \leq \mathcal{R}_{\pm}, r \geq 1/4 \mid \mathcal{R}_+ - \mathcal{R}_- \mid$ [2].

Decay modes with excited kaons, such as $B^- \rightarrow DK^{*-}$ can be used in a similar way to constrain ϕ_3 . Belle studied these channels with 88 fb^{-1} of data [7]. K^{*-} was reconstructed from $K^{*-} \rightarrow K_s^0 \pi^-, K_s^0 \rightarrow \pi^+ \pi^-$ decays. The measured branching fraction $\mathcal{B} = (5.2 \pm 0.5(\text{stat}) \pm 0.6(\text{sys})) \times 10^{-4}$ is consistent with earlier CLEO result based on much lower statistics [6]. The performed analysis showed no indication of a non-resonant $D^0 K_s^0 \pi^-$ contribution. Event yields obtained in bins of $K_s^0 \pi^-$ invariant mass and K^{*-} helicity angle are shown in Fig. 4 a and 4 b respectively. Results are consistent with Monte Carlo simulations of pure $B^- \rightarrow D^0 K^{*-}$ decay.

With this data sample Belle observed for the first time $B^- \rightarrow D_{\pm} K^{*-}$ decays (see Fig. 5). The significance is 4.3σ for CP-even ($K^+ K^-, \pi^+ \pi^-$) and 2.4σ for CP-odd ($K_s^0 \pi^0, K_s^0 \phi, K_s^0 \omega$) states. Event yields are substantially lower than for $B^- \rightarrow D_{\pm} K^-$, which is due to the limited reconstruction efficiency of K^{*-} . Preliminary measurements of CP asymmetries are presented in Table 2.

2.2 $B^- \rightarrow D^* K^{*-}$

Decay modes into two vector mesons can be also used to measure ϕ_3 [9]. An interesting feature of these channels

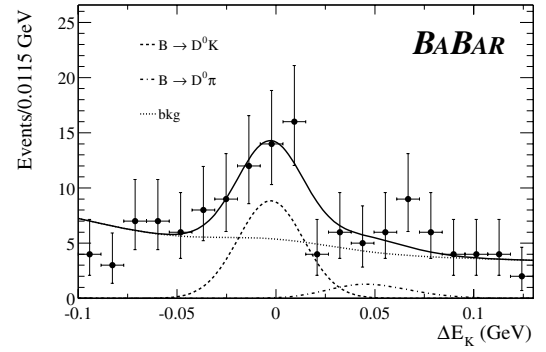


Fig. 2. ΔE distributions of $B^{\pm} \rightarrow D_{\pm} K^{\pm}$ from BaBar. The kaon mass is assumed for the prompt hadron. Contributions from $B^{\pm} \rightarrow D_{\pm} \pi^{\pm}$ are seen at $\Delta E \approx 0.05 \text{ GeV}$

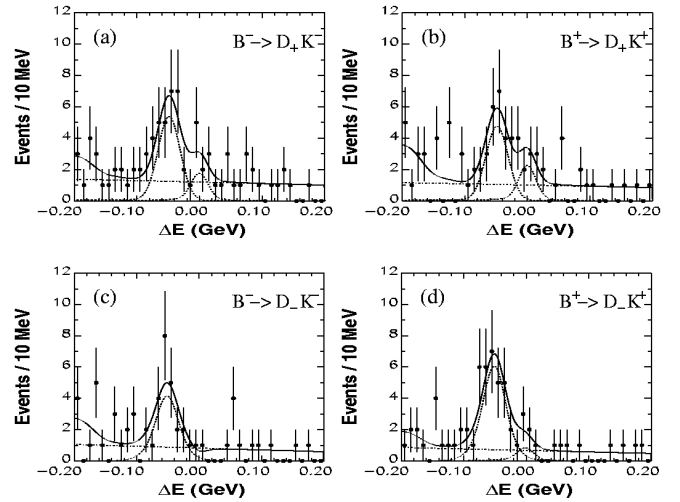


Fig. 3. ΔE distributions of $B^{\pm} \rightarrow D_{\pm} K^{\pm}$ from Belle. The pion mass is assumed for the prompt hadron. Contributions from $B^{\pm} \rightarrow D_{\pm} \pi^{\pm}$ are seen at $\Delta E = 0$

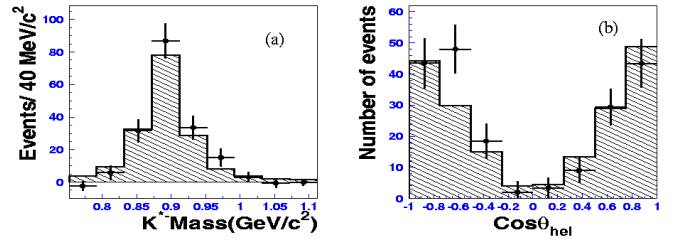


Fig. 4. Yields for $B^{\pm} \rightarrow D^0 K^{\pm}$ from Belle in bins of $K_s^0 \pi^-$ invariant mass (a) and in the bins of $\cos \theta_{hel}$ (b). The hatched histograms are Monte Carlo simulations of $B^- \rightarrow D^0 K^{*-}$

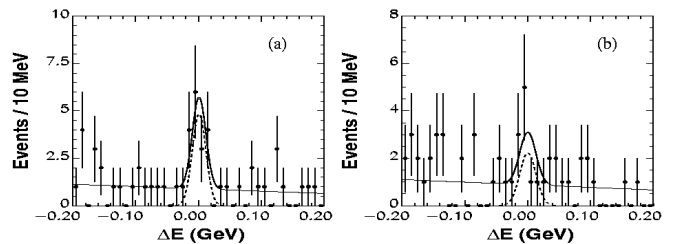


Fig. 5. ΔE distributions of $B^{\pm} \rightarrow D_{\pm} K^{\pm}$ (a) and $B^{\pm} \rightarrow D_{\pm} K^{*\pm}$ (b) from Belle

Table 1. Measurements of $R = \frac{\mathcal{B}(B^- \rightarrow D^0 K^-)}{\mathcal{B}(B^- \rightarrow D^0 \pi^-)}$. The first error is statistical and the second is systematic

experiment	R	data sample
CLEO [3]	$(9.9 \pm 1.3 \pm 0.7)\%$	15.3 fb^{-1}
BaBar [4]	$(8.31 \pm 0.25 \pm 0.2)\%$	56 fb^{-1}
Belle [5]	$(7.7 \pm 0.5 \pm 0.6)\%$	79 fb^{-1}

Table 2. Results on \mathcal{A}_\pm and \mathcal{R}_\pm . The first error is statistical and the second is systematic

mode	\mathcal{A}_\pm	\mathcal{R}_\pm	exp.
$D_+ K^-$	$0.17 \pm 0.23 \pm 0.08$	$1.06 \pm 0.26 \pm 0.17$	BaBar
$D_+ K^-$	$0.06 \pm 0.19 \pm 0.04$	$1.21 \pm 0.25 \pm 0.14$	Belle
$D_- K^-$	$-0.19 \pm 0.17 \pm 0.05$	$1.41 \pm 0.27 \pm 0.15$	Belle
$D_+ K^{*-}$	$-0.02 \pm 0.33 \pm 0.07$	-	Belle
$D_- K^{*-}$	$0.19 \pm 0.50 \pm 0.04$	-	Belle

is the presence of additional observables resulting from interference between helicity amplitudes.

BaBar has recently released results on the $B^- \rightarrow D^{*0} K^{*-}$ decays based on data sample of 79 fb^{-1} [8]. The measured branching fraction of $(8.3 \pm 1.1(\text{stat}) \pm 1.0(\text{sys})) \times 10^{-4}$ agrees well with the earlier CLEO result [6]. BaBar also presented the first measurement of the fraction of longitudinal polarization $\Gamma_L/\Gamma = 0.86 \pm 0.06(\text{stat}) \pm 0.03(\text{sys})$ in these decays.

2.3 $\bar{B}^0 \rightarrow D\bar{K}^{(*)0}$

Neutral B decays of the type (1) are of special interest. In this case the amplitude $A_{b \rightarrow c}$ is also color suppressed. This leads to decay rates as low as $\approx 10^{-5}$, but at the same time the larger amplitude ratio $r \sim 0.4$ augments interference effects ([10]). The flavor of B^0 can be tagged in the modes with excited kaons from $\bar{K}^{*0} \rightarrow K^- \pi^+$ decay products. The ϕ_3 can be constrained from decays of this type in a similar way as from charged B decays [11]. When the final states contain K_s^0 , $B^0 - \bar{B}^0$ mixing leads to time-dependent asymmetries, which are sensitive to a combination of mixing and decay "phases" ([10], [12])." Belle, using a data sample of 78 fb^{-1} , observed for the first time the decays $\bar{B}^0 \rightarrow D^0 \bar{K}^0$ and $\bar{B}^0 \rightarrow D^0 \bar{K}^{*0}$ (Fig. 6) with the branching fractions $\mathcal{B}(\bar{B}^0 \rightarrow D^0 \bar{K}^0) = (5.0_{-1.2}^{+1.3}(\text{stat}) \pm 0.6(\text{sys})) \times 10^{-5}$ and $\mathcal{B}(\bar{B}^0 \rightarrow D^0 \bar{K}^{*0}) = (4.8_{-1.0}^{+1.1}(\text{stat}) \pm 0.5(\text{sys})) \times 10^{-5}$ [13]. Also in this case no indication of non-resonant contribution to $\bar{B}^0 \rightarrow D^0 \bar{K}^{*0}$ was found.

3 Summary

Belle and BaBar continue a steady progress in the studies of the decays of the type (1). Among the most important results one should mention measurements of charged B decays to D CP-eigenstates which provide the complete

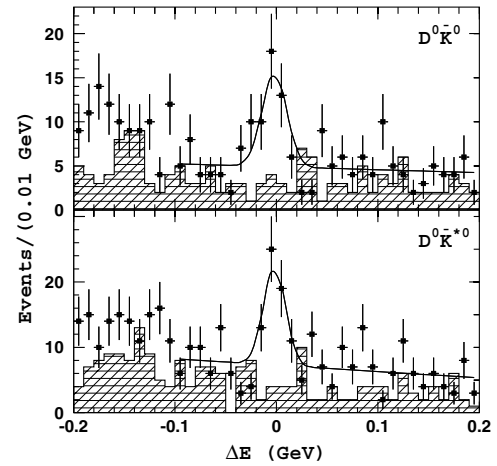


Fig. 6. ΔE distributions for $\bar{B}^0 \rightarrow D^0 \bar{K}^{(*)0}$ from Belle. The hatched histograms show the distributions of events in D^0 side-band regions

set of observables needed in the GLW method and observation of the color suppressed $\bar{B}^0 \rightarrow D^0 \bar{K}^{(*)0}$ decays. With new theoretical approaches and the excellent performance of KEKB and PEP-II, one can hope for meaningful constraints on ϕ_3 in the near future.² This also gives good prospects for ϕ_3 measurements at B-factories with upgraded luminosities.

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² After this conference Belle presented 90% confidence interval for the ϕ_3 : $61^\circ \leq \phi_3 \leq 142^\circ$ [14]. The result was obtained with 140 fb^{-1} of data using a novel method which employs interference between three-body D^0 and \bar{D}^0 decays [15].